

SOLUTION TO NONLINEAR GRADIENT DEPENDENT SYSTEMS WITH A BALANCE LAW

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ABSTRACT . In this paper , we are concerned with the initial boundary value problem (IBVP) and with the Cauchy problem to the reaction - diffusion system

$$\begin{aligned} u_t - \Delta u &= -u^n |\nabla v|^p, \\ v_t - d\Delta v &= u^n |\nabla v|^p, \end{aligned}$$

where $1 \leq p \leq 2$, d and n are positive real numbers . Results on the existence and large - time behavior of the solutions are presented .

1 . INTRODUCTION

In the first part of this article , we are interested in the existence of global classical nonnegative solutions to the reaction - diffusion equations

$$u_t - \Delta_v u = -u^n d \frac{|\nabla v|^p}{|\nabla v|} =: -f(u, v), \tag{1.1}$$

posed on $\mathbb{R}^+ \times \Omega$ with initial data

$$u(0; x) = u_0(x), \quad v(0; x) = v_0(x) \quad \text{in } \Omega \tag{1.2}$$

and boundary conditions (in the case Ω is a bounded domain in \mathbb{R}^n)

$$\frac{\partial u}{\partial \eta} = \frac{\partial v}{\partial \eta} = 0, \quad \text{on } \mathbb{R}^+ \times \partial\Omega. \tag{1.3}$$

Here Δ is the Laplacian operator , u_0 and v_0 are given bounded nonnegative functions , $\Omega \subset \mathbb{R}^n$ is a regular domain , η is the outward normal to $\partial\Omega$. The diffusive coefficient d is a positive real . One of the basic questions for (1 . 1) - (1 . 2) or (1 . 1) - (1 . 3) is the existence of global solutions . Motivated by extending known results on reaction - diffusion systems with conservation of the total mass but with nonlinearities depending only for the unknowns , Boudiba , Mouley and Pierre succeeded in obtaining L^1 solutions only for the case $u^n |\nabla v|^p$ with $p < 2$. In this article , we are interested essentially in classical solutions in the case where $p = 2$ (Ω bounded or $\Omega = \mathbb{R}^n$; in the latter case , there are no boundary conditions) .

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2. RESULTS

The existence of a unique classical solution over the whole time interval $[0, T_{\max}]$ can be obtained by a known procedure : a local solution is continued globally by using a priori estimates on $\|u\|_{\infty}, \|v\|_{\infty}, \|\nabla u\|_{\infty}$, and $\|\nabla v\|_{\infty}$. **The Cauchy problem . Uniform bounds for u and v .** First, we consider the auxiliary problem

$$L_{\lambda}\omega := \omega_{t\omega}^{-\lambda\Delta\omega} = b\nabla\omega\omega_{(x)} \in tL_{\infty}^{>0}, x \in \mathbb{R}^N \quad (2.1)$$

where $b = (b_1(t, x), \dots, b_N(t, x)), b_i(t, x)$ are continuous on $[0, \infty) \times \mathbb{R}^N$, ω is a classical solution of (2.1). **Lemma 2.1.** Assume that $\omega_t, \nabla\omega, \omega_{x_i x_i}, i = 1, \dots, N$ are continuous,

$$L_{\lambda}\omega \leq 0, \quad (\geq) \quad (0, \infty) \times \mathbb{R}^N \quad (2.2)$$

and $\omega(t, x)$ satisfies (2.1) 2. Then

$$\begin{aligned} \omega(t, x) &\leq C := \sup_{x \in \mathbb{R}^N} \omega_0(x), \quad (0, \infty) \times \mathbb{R}^N. \\ \omega(t, x) &\geq C := \inf_{x \in \mathbb{R}^N} \omega_0(x), \quad (0, \infty) \times \mathbb{R}^N. \end{aligned}$$

The proof of the above lemma is elementary and hence is omitted. Now, we consider the problem (1.1) - (1.2). It follows by the maximum principle that

$$u, v \geq 0, \quad \text{in } \mathbb{R}^+ \times \mathbb{R}^N.$$

Uniform bounds of u . We have

$$u \leq C_1 := \sup_{\mathbb{R}^N} u_0(x),$$

thanks to the maximum principle. Uniform bounds of v . Next, we derive an upper estimate for v . Assume that $1 \leq p < 2$. We transform (1.1) 2 by the substitution $\omega = e^{\lambda v} - 1$ into

$$\omega_t - \lambda\Delta\omega = \lambda e^{\lambda v}(v_t - d\Delta v - d\lambda |\nabla v|^2) = \lambda e^{\lambda v}(u^n |\nabla v|^p - d\lambda |\nabla v|^2).$$

Let

$$\phi(x) \equiv Cx^p - d\lambda x^2; \quad C > 0, x \geq 0.$$

By elementary computations,

$$\phi(x) \geq 0 \quad \text{when } x \leq \left(\frac{C}{\lambda d}\right)^{1/(2-p)}.$$

But in this case

$$|\nabla v| \leq \left(\frac{c}{\lambda d}\right)^{1/(2-p)}.$$

In the case $x \geq \left(\frac{c}{\lambda d}\right)^{1/(2-p)}$,

$$\phi(x) \leq 0 \tag{2.3}$$

and hence $\omega \leq M$ where

$$M = C \left(\frac{pC}{2d\lambda} \right)^{p/2-p} \left(\frac{2-p}{2} \right). \tag{2.4}$$

Then we have $v \leq C_2$.

$$L_d v + kv = kv + u^n |\nabla v|^p \tag{2.5}$$

and transform it by the substitutions $\omega = e^{kt}v$ to obtain

$$L_d \omega = e^{kt}(L_d v + kv) = e^{kt}(kv + u^n |\nabla v|^p), \quad t > 0, x \in \mathbb{R}^N$$

$$\omega(0, x) = v_0(x).$$

Now let

$$G_\lambda = G_\lambda(t - \tau; x - \xi) = \frac{1}{[4\pi\lambda(t - \tau)]^{\frac{N}{2}}} \exp\left(-\frac{|x - \xi|^2}{4\lambda(t - \tau)}\right)$$

be the fundamental solution related to the operator L_λ . Then, with $Q_t = (0, t) \times \mathbb{R}^N$, we have

$$\omega = e^{kt}v = v^0(t, x) + \int_{Q_t} G_d(t - \tau; x - \xi) e^{k\tau} (kv + u^n |\nabla v|^p) d\xi d\tau$$

or

$$v = e^{-kt}v^0 + \int_{Q_t} e^{-k(t-\tau)} G_d(t - \tau; x - \xi) (kv + u^n |\nabla v|^p) d\xi d\tau, \tag{2.6}$$

where $v^0(t, x)$ is the solution of the homogeneous problem

$$L_d v^0 = 0, \quad v^0(0, x) = v_0(x).$$

From (2.6) we have

$$|\nabla v| = e^{-kt} |\nabla v^0| + \int_{Q_t} e^{-k(t-\tau)} |\nabla_x G_d(t - \tau; x - \xi)| (kv + u^n |\nabla v|^p) d\xi d\tau. \tag{2.7}$$

Now we set $\nu_1 = \sup |\nabla v|$ and $1_\nu^0 = \sup |\nabla v^0|$, in Q_t . From (2.7), and using $v \leq C_2$, we have

$$\nu_1 = 1_\nu^0 + (kC_2 + C_1^n \nu_1^p) \int_0^t e^{-k(t-\tau)} \left(\int_{\mathbb{R}^N} |\nabla_x G_d(t - \tau; x - \xi)| d\xi \right) d\tau.$$

We also have

$$\int_{\mathbb{R}^N} |\nabla_x G_d(t - \tau; x - \xi)| d\xi = \int_{\mathbb{R}^N} \frac{|x - \xi|}{2d(t - \tau)} |G_d(t - \tau; x - \xi)| d\xi$$

which is transformed by the substitution $\rho = 2\sqrt{d(t - \tau)}\nu$ into

$$\int_{\mathbb{R}^N} |\nabla_x G_d| d\rho = \frac{w_N}{\pi^{N/2}} \int_0^\infty e^{-\nu^2} d\nu = \frac{\chi}{\sqrt{d(t - \tau)}}$$

where $\chi = \text{line - slash}_{2\pi} w_N N^N 2\Gamma(\frac{N+1}{2}) = \frac{\Gamma(\frac{N+1}{2})}{\Gamma(\frac{N}{2})}$. It follows that

$$\nu_1 = 1_\nu^0 + (kC_2 + C_1^n \nu_1^p) \frac{\chi}{\sqrt{d}} \int_0^t e^{-k(t-\tau)} \frac{d\tau}{\sqrt{t - \tau}}. \tag{2.8}$$

Recall that

$$\int_0^t e^{-k(t-\tau)} \frac{d\tau}{\sqrt{t-\tau}} = \frac{2}{\sqrt{k}} \int_0^t e^{-z^2} dz < \sqrt{\frac{\pi}{k}}.$$

If we set $s = \sqrt{k}$ in (2 . 8) then we have

$$\nu_1 \leq 1_\nu^0 + (sC_2 + \frac{C_1^n}{s} \nu_1^p) \chi \sqrt{\frac{\pi}{d}}. \quad (2.9)$$

4 Z. DAHMANI, S. KERBAL EJDE - 2017 / 158 Now we minimize the right hand side of (2.9) with respect to s to obtain

$$\nu_1 \leq 1_\nu^0 + \frac{2\chi\sqrt{\pi}}{d}(C_2 C_1^n p_{\nu_1})^{1/p}. \quad (2.10)$$

Notethat $1_\nu^0 = C_2$.

We have two cases : Case (i) $1 \leq p < 2$. In this case (2.10) implies

$$|\nabla v| \leq \nu_1 \leq \dots \nu(p) = D, \quad \text{in } Q_t, \quad (2.11)$$

where D is a positive constant .

Case (ii) $p = 2$. In this case (2.10) holds under the additional condition

$$C_2 C_1^n \leq \frac{d}{4\pi\chi}. \quad (2.12)$$

Similarly we obtain from (1.1) ,

$$U_1 := \sup_{Q^T} |\nabla u| \leq C_1 + C_1 \frac{2\sqrt{\pi}\chi}{\sqrt{d}} 1_\nu^{p/2} \leq \text{Constant}. \quad (2.13)$$

The estimates (2.10) and (2.13) are independent of t , hence $T_{\max} = +\infty$.

Finally , we have the main result .

Theorem 2.2 . *Let $p = 2$ and (u_0, v_0) be bounded such that (2.12) holds then system (1.1) - (1.2) admits a global solution .*

2.2 . The Neumann Problem . In this section , we are concerned with the Neumann problem

$$u_{v_t} - \Delta_{d\Delta v} u = -u^n \frac{|\nabla v|^2}{|\nabla v|^2} \quad (2.14)$$

where Ω be a bounded domain in \mathbb{R}^N , with the homogeneous Neumann boundary condition

$$\frac{\partial u}{\partial \nu} = \frac{\partial v}{\partial \nu} = 0, \quad \text{on } \mathbb{R}^+ \times \partial\Omega \quad (2.15)$$

subject to the initial conditions

$$u(0; x) = u_0(x); \quad v(0; x) = v_0(x) \quad \text{in } \Omega. \quad (2.16)$$

The initial nonnegative functions u_0, v_0 are assumed to belong to the Holder space

$$C^{2,\alpha}(\Omega).$$

Uniform bounds for u and v . In this section a priori estimates on $\|u\|_\infty$ and $\|v\|_\infty$ are presented . **Lemma 2.3 .** *For each $0 < t < T_{\max}$ we have*

$$0 \leq u(t, x) \leq M, \quad 0 \leq v(t, x) \leq M, \quad (2.17)$$

for any $x \in \Omega$.

Proof . Since $u_0(x) \geq 0$ and $f(0, v) = 0$, we first obtain $u \geq 0$ and then $v \geq 0$ as $v_0(x) \geq 0$. Using the maximum principle , we conclude that

$$0 \leq u(t, x) \leq M, \quad \text{on } QT$$

where

$$M \geq M_1 := \max_{x \in \Omega} u_0(x).$$

$$\begin{aligned} \omega_t - d\Delta\omega &= \lambda |\nabla v|^2 (u^n - d\lambda)e^{\lambda v}, \quad \text{on } QT \\ \frac{\partial u}{\partial \nu} &= 0 \quad \text{on } \partial S_T. \end{aligned}$$

Consequently as $d\lambda > \max_{\Omega} u^n$, we deduce from the maximum principle that

$$0 \leq \omega(t, x) \leq \exp(\lambda \|v_0\|_{\infty}) - 1.$$

Hence

$$v(x, t) \leq \frac{1}{\lambda} \ln(\|\omega\|_{\infty} + 1) \leq \text{Constant} < \infty.$$

□

Uniform bounds for $|\nabla v|$ and $|\nabla u|$. To obtain uniform a priori estimates for $|\nabla v|$, we make use of some techniques already used by Tomi [8] and von Wahl [9]

Lemma 2.4. *Let (u, v) be a solution to (2.10) - (2.12) in its maximal interval of existence $[0, T_{\max}[$. Then there exist a constant C such that*

$$\|u\|_{L^{\infty}([0, T[, W^2, q(\Omega))} \leq C \quad \text{and} \quad \|v\|_{L^{\infty}([0, T[, W^2, q(\Omega))} \leq C.$$

Proof. Let us introduce the function

$$f_{\sigma, \epsilon}(t, x, u, \nabla v) = \sigma u^n(t, x) \frac{\epsilon + |\nabla v|^2}{1 + \epsilon |\nabla v|^2}.$$

It is clear that $|f_{\sigma, \epsilon}(t, x, u, \nabla v)| \leq C(1 + |\nabla v|^2)$ and a global solution $v_{\sigma, \epsilon}$ differentiable in σ for the equation

$$v_t - d\Delta v = f_{\sigma, \epsilon}(t, x, u, \nabla v)$$

exists. Moreover, $v_{\sigma, \epsilon} \rightarrow v$ as $\sigma \rightarrow 1$ and $\epsilon \rightarrow 0$, uniformly on every compact of

$$[0, T_{\max}[.$$

The function $\omega_{\sigma} := \frac{\partial v_{\sigma, \epsilon}}{\partial \sigma}$ satisfies

$$\partial_t \omega_{\sigma} - d\Delta \omega_{\sigma} = u^n(t, x) \frac{\epsilon + |\nabla v_{\sigma}|^2}{1 + \epsilon |\nabla v_{\sigma}|^2} - 2\sigma u^{n(\epsilon)} \frac{2 - 1) \nabla v_{\sigma} \cdot \nabla \omega_{\sigma}}{(1 + \epsilon |\nabla v_{\sigma}|^2)^2}. \quad (2.18)$$

Hereafter, we derive uniform estimates in σ and ϵ . Using Solonnikov's estimates for parabolic equation [5] we have

$$\|\omega_{\sigma}\|_{L^{\infty}([0, T(u_0, v_0)[, W^2, p(\Omega))} \leq C[\|\nabla v_{\sigma}\|_{L^p(\Omega)}^2 + \|\nabla v_{\sigma} \cdot \nabla \omega_{\sigma}\|_{L^p(\Omega)}^2].$$

The Gagliardo - Nirenberg inequality [5] in the form

$$\|u\|_{W^1, 2p(\Omega)} \leq C \|u\|_{L^{\infty}(\Omega)}^{1/2} C \|u\|_{W^2, p(\Omega)}^{1/2}$$

and the δ -Young inequality (where $\delta > 0$)

$$\alpha\beta \leq \frac{1}{2}\left(\delta\alpha^2 + \frac{\beta^2}{\delta}\right),$$

allows one to obtain the estimate

$$\|\omega_\sigma\|_{L^\infty([0, T(u_0, v_0)[, W^2, p(\Omega))} \leq C(1 + \|\omega_\sigma\|_{W^2, p(\Omega)}).$$

But $\omega_\sigma = \frac{\partial v_\sigma}{\partial \sigma}$, hence by Gronwall ' s inequality we have

$$\|v_\sigma\|_{L^\infty([0, T[, W^2, p(\Omega))} \leq Ce^{C\sigma}.$$

$$\|v\|_{L^\infty([0, T], W^{2, p}(\Omega))} \leq C.$$

On the other hand, the Sobolev injection theorem allows to assert that $u \in C^{1, \alpha}(\Omega)$. Hence in particular $|\nabla u| \in C^{0, \alpha}(\Omega)$. Since $|\nabla v|$ is uniformly bounded, it is easy then to bound $|\nabla u|$ in $L^\infty(\Omega)$. As a consequence, one can affirm that the solution (u, v) to problem (2.14) - (2.16) is global; that is $T_{\max} = \infty$. \square

2.3. Large-time behavior. In this section, the large time behavior of the global solutions to (2.14) - (2.16) is briefly presented.

Theorem 2.5. *Let $(u_0, v_0) \in C^{2, \epsilon}(\Omega) \times C^{2, \epsilon}(\Omega)$ for some $0 < \epsilon < 1$. The system (2.14) - (2.16) has a global classical solution. Moreover, as $t \rightarrow \infty$, $u \rightarrow k_1$ and $v \rightarrow k_2$ uniformly in x , and*

$$k_1 + k_2 = \frac{1}{|\Omega|} \int_{\Omega} [u_0(x) + v_0(x)] dx.$$

Proof. The proof of the first part of the Theorem is presented above. Concerning the large time behavior, observe first that for any $t \geq 0$,

$$\int_{\Omega} [u(t, x) + v(t, x)] dx = \int_{\Omega} [u_0(x) + v_0(x)] dx.$$

Then, the function $t \rightarrow \int_{\Omega} u(x) dx$ is bounded; as it is decreasing, we have

$$\int_{\Omega} u(x) dx \rightarrow k_1 \text{ as } t \rightarrow \infty;$$

the function $t \rightarrow \int_{\Omega} v(x) dx$ is increasing and bounded, hence admits a finite limit k_2 as $t \rightarrow \infty$. As $\bigcup_{t \geq 0} \{(u(t), v(t))\}$ is relatively compact in $C(\text{---}\Omega) \times C(\text{---}\Omega)$,

$$u(\tau_n) \rightarrow \tilde{u}, \quad v(\tau_n) \rightarrow \tilde{v} \text{ in } C(\text{---}\Omega),$$

through a sequence $\tau_n \rightarrow \infty$. It is not difficult to show that in fact (\tilde{u}, \tilde{v}) is the stationary solution to (2.14) - (2.16) (see [3]).

As the stationary solution (u_s, v_s) to (2.14) - (2.16) satisfies

$$\begin{aligned} -\Delta u_s &= -u_s^n |\nabla v_s|^2, \quad \text{in } \Omega, \\ -d\Delta v_s &= u_s^n |\nabla v_s|^2, \quad \text{in } \Omega, \quad \frac{\partial u_s}{\partial \nu} = \frac{\partial v_s}{\partial \nu} = 0, \quad \text{on } \partial\Omega, \end{aligned}$$

we have

$$-\int_{\Omega} \Delta u_s \cdot u_s dx = -\int_{\Omega} u_s^{n+1} |\nabla v_s|^2 dx$$

which in the light of the Green formula can be written

$$\int_{\Omega} |\nabla u_s|^2 dx = -\int_{\Omega} u_s^{n+1} |\nabla v_s|^2 dx$$

hence $|\nabla u_s| = |\nabla v_s| = 0$ implies $u_s = k_1$ and $v_s = k_2$. \square

Remarks. (1) It is very interesting to address the question of existence global solutions of the system (2.14) - (2.16) with a genuine nonlinearity of the form $u^n |\nabla v|^p$ with $p \geq 2$.

(2) It is possible to extend the results presented here for systems with nonlinear boundary conditions satisfying reasonable growth restrictions.

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